

Modeling the wake of the Marquesas archipelago

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Key Points:

- Characterizing the eddy activity in the island wake of the Marquesas archipelago using high-resolution ocean model
- Revealing the importance of the fine-scale dynamics within the Island Mass Effect

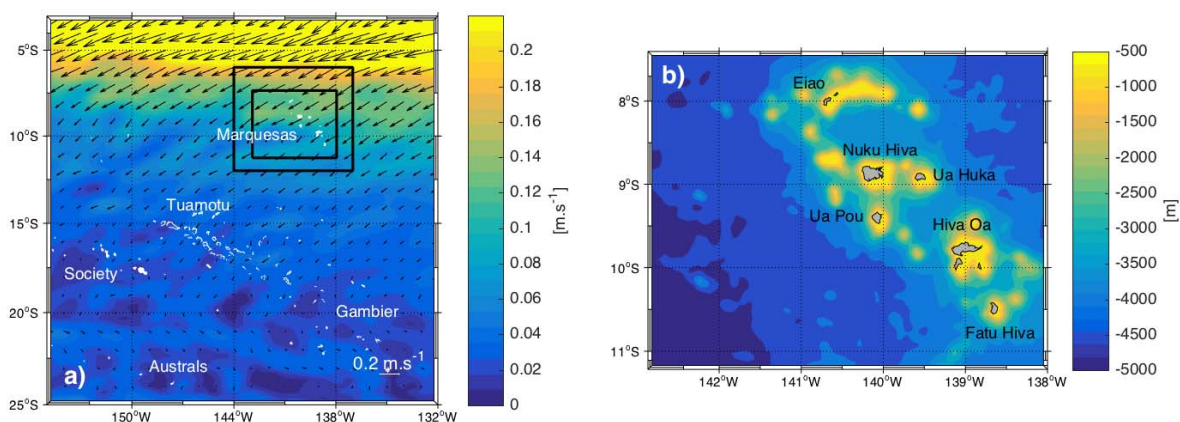
24 **Abstract**

25 In this study, a high-resolution (~ 2.5 km) numerical model was set up to investigate the fine-
 26 scale activity within the region of the Marquesas archipelago. This has never been performed
 27 before. The robustness of the model results is assessed by comparison with remote sensing and in
 28 situ observations. Our results highlight regions of warm waters leeward of the different islands
 29 with high eddy kinetic energy (*EKE*) on their sides. The analysis of energy conversion terms
 30 reveals contributions to *EKE* variability by wind, baroclinic and barotropic instabilities. The use
 31 of a geometry-based eddy detection algorithm reveals the generation of cyclonic and anticyclonic
 32 eddies in the wake of the largest islands, with both an inshore and offshore effect. Maximum
 33 eddy activity occurs in austral winter following the seasonality of both wind stress and *EKE*
 34 intensity. Most eddies have a radius between 20 and 30 km and are generally cyclonic rather than
 35 anticyclonic. Significant vertical velocities are observed in the proximity of the islands,
 36 associated with topographically-induced flow separation. Eddy trapping inshore waters are
 37 advected offshore in the wake of the islands. The overall influence of these fine-scale dynamics
 38 could explain the strong biological enhancement of the archipelago.

39 **1 Introduction**

40 The Marquesas archipelago (144°W - 137°W ; 8°S - 11°S) is located in the northern part of French
 41 Polynesia, central South Pacific, where the South Equatorial Current (SEC) flows south-
 42 westward (Figure 1a). It is composed by a dozen of small volcanic islands with mountains up to
 43 1224 m, rugged steep cliffs and no surrounding coral reefs (*Savanier et al.*, 2006; Figure 1b).
 44 The main five islands are Nuku Hiva (339 km^2), Ua Pou (105 km^2) and Ua Huka (83 km^2) in the
 45 northern part of the archipelago, and Hiva Oa (320 km^2) and Fatu Hiva (85 km^2) in the southern
 46 part. Despite their relatively small area coverage, a remarkable plume of chlorophyll-a (Chl, a
 47 proxy of phytoplankton biomass) can be observed leeward of the islands in the open ocean from
 48 satellite derived ocean color [*Signorini et al.*, 1999; *Martinez and Maamaatuaiahutapu*, 2004].
 49 Such biological enhancement is referred to as an island mass effect (IME) [*Doty, M. S., & Oguri*,
 50 1956]. While several physical processes can be involved in the present IME, such as coastal
 51 upwelling, Ekman pumping, eddies or internal waves [*Heywood et al.*, 1990, 1996; *Barton*,
 52 2001; *Sangrà et al.*, 2001; *Palacios*, 2002], the IME dominant mechanisms have not been
 53 elucidated yet for the Marquesas archipelago.

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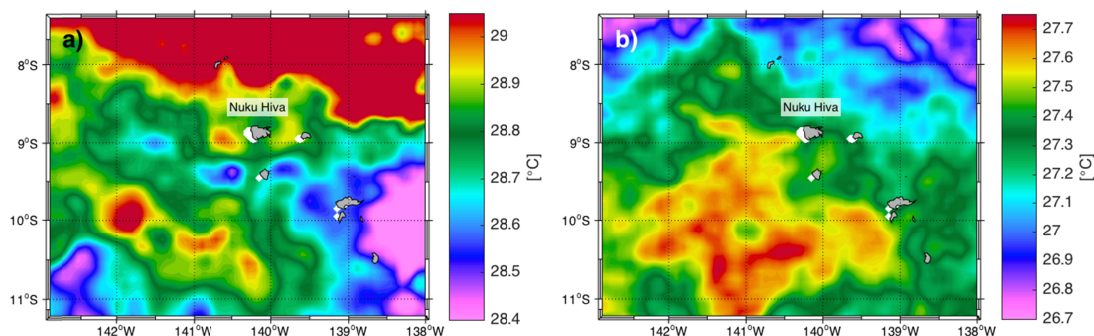


55

56 Figure 1: (a) Mean surface current from the satellite derived Ocean Surface Current Analysis –
57 Real time (OSCAR) product (time averaged over Oct-1992 to Jun-2015, in m s^{-1}) over French
58 Polynesia. The islands are represented in white, as well as the name of the five archipelagos. The
59 black boxes show the parent and child grids implemented in the ROMS configuration. (b)
60 Bottom topography from the 2-arc minute topography/bathymetry dataset ETOPO2 (in m)
61 around the archipelago used in the model configuration as well as the names of the main islands.

62
63 It is well known that oceanic currents encountering islands generate complex wakes. Several
64 studies in other regions have assessed these island wakes and their forcing mechanisms using a
65 remote-sensing [Wolanksi *et al.*, 1996; DiGiacomo and Holt, 2001], in situ data [Hasegawa *et al.*,
66 2004, 2008] or a modeling approach [Dietrich *et al.*, 1996; Wolanksi *et al.*, 1996; Dong *et al.*,
67 2007; Jiménez *et al.*, 2008; Hasegawa *et al.*, 2009]. An important distinction has to be made
68 between shallow- and deep-water wakes. Shallow-water wakes occur when considering islands
69 in shallow shelves or estuaries where near-shore bottom drag acts as the primary source of
70 vorticity generation [Wolanksi *et al.*, 1996; Alaei *et al.*, 2004; Neill and Elliott, 2004]. Deep-
71 water wakes occur leeward tall islands surrounding by a deep bathymetry - like in the Marquesas
72 archipelago - where bottom influence can be neglected and topographic and wind forcing are the
73 primary sources of vorticity generation. Topographic forcing refers to the detachment of the
74 frictional layer around the island as a result of differential bottom stress imparted to the flow by
75 the sloping sides of the island [Dong *et al.*, 2007]. Wind forcing refers to eddy generation as a
76 consequence of Ekman pumping induced by wind shear in the island's wake [Jiménez *et al.*,
77 2008].

78 Eddy formation and propagation in island wakes could play a key role in the IME of the
79 Marquesas archipelago. Indeed, a wind curl dipole has been reported in the archipelago
80 [Martinez *et al.*, 2009] possibly forming oceanic cyclonic eddies and jets inducing an enrichment
81 of the upper layer leeward the islands by mechanisms such as the ones proposed by Hasegawa *et al.*
82 [2009] and Andrade *et al.* [2014]. Nevertheless, the previous studies on the Marquesas IME
83 using satellite altimetry did not mention any eddy features [Signorini *et al.*, 1999; Legeckis *et al.*,
84 2004; Martinez and Maamaatuaiahutapu, 2004]. This is probably due to the coarseness of the
85 Sea Level Anomaly (SLA) products used in these studies. However, some eddy activity does
86 imprint the Sea Surface Temperature (SST). For example, on Nov. 17, 2015 (Figure 2a) a SST
87 dipole is observed leeward Nuku Hiva at (140.5W; 9°S), as well as a warm eddy further
88 southwest, while a cool eddy is present at (141°W; 10.1°S) on Aug. 21, 2012 (Figure 2b). In both
89 cases, it is important to note that the size of the structures is of the same order than the island
90 diameter (~30 km for Nuku Hiva). Such an eddy activity could possibly be at the origin of
91 nutrient uplift from deep rich waters to the euphotic zone allowing the development of
92 phytoplankton. Because the spatio-temporal resolution of satellite altimetry is too coarse to
93 investigate the small scales of such aforementioned eddy activity around the Marquesas
94 archipelago, recourse to high-resolution modeling taking into account the topographic forcing
95 (i.e., the presence of the islands) has been done by the present study. Documenting the
96 generation and characteristics of the oceanic eddies should be viewed as a first step toward an
97 explanation of the origin of the surface layer biological enrichment in this archipelago.



98
99 Figure 2: SST ($^{\circ}\text{C}$) from Geostationary Operational Environmental Satellite (GOES)/ Polar
100 Operational Environmental Satellite (POES) in the Marquesas archipelago for (a) 2015-NOV-17
101 and (b) 2012-AUG-21.
102

103 The present work is organized as follows: Section 2 presents the data and the methodology. A
104 comparison of the oceanic conditions between the model, satellite and *in situ* observations is
105 performed in section 3. In this section, eddy properties generated leeward the islands are also
106 characterized. Finally, section 4 concludes this study.

107 2 Data and Methods

108 2.1 Numerical model

109 Our model is based on the Regional Ocean Modeling System (ROMS)-AGRIF (Adaptive Grid
110 Refinement in Fortran) code provided by *Debreu et al.* [2012] and *Penven et al.* [2007]. The
111 ROMS model is a split-explicit, free surface and terrain-following vertical coordinate oceanic
112 model [*Shchepetkin and McWilliams*, 2003; 2005]. The AGRIF version is especially suitable to
113 study regional scale since it has the ability to manage an arbitrary number of fixed grids and
114 embedding levels. We configure the two-way embedding procedure. It means that the parent grid
115 provides the boundary conditions for the child grid, and that the solution of the child grid is used
116 to improve the larger scale parent grid solution, allowing a smooth, continuous interfacing
117 between grid levels [*Debreu and Blayo*, 2008; *Debreu et al.*, 2012]. Therefore, we defined two
118 embedded grids: 1) the parent grid extends from 137°W to 144°W and 6°S to 12°S ; 2) the child
119 grid extends from 138°W to 143°W and 7.3°S to 11.3°S as shown in Figure 1a. The grid
120 refinement rate is 3, implying a $1/15^{\circ}$ ($\sim 7\text{km}$) and $1/45^{\circ}$ ($\sim 2.5\text{ km}$) grid resolutions for the parent
121 and child grids, respectively. This allows a sufficient sampling of the island topography and of
122 the fine-scale dynamics, as the first baroclinic Rossby radius of deformation is about 120 km in
123 this region according to *Chelton et al.* [1998]. The internal (external) time stepping is set to 3600
124 (600) sec for the parent grid and 1200 (200) sec for the child grid. Both grids have 32 vertical
125 levels and the vertical s-coordinate is stretched for boundary layer resolution. The topography is
126 derived from the $2'$ resolution ETOPO2 database provided by NOAA-NGDC [*Smith and*
127 *Sandwell*, 1997]. The bathymetry field has been filtered to keep the slope parameter < 0.25
128 [*Beckmann and Haidvogel*, 1993]. The K-profile parameterization (KPP) vertical mixing scheme
129 from *Large et al.* [1994] is used to parameterize vertical mixing processes.

130 All the external forcings of the ROMS simulations are based on monthly climatologies. At the
131 surface, the heat and fresh water fluxes are extracted on a monthly $\frac{1}{2}^{\circ}$ grid from the

132 Comprehensive Ocean-Atmosphere dataset (COADS) [da Silva et al., 1994]. The wind forcing is
 133 issued from the QuikSCAT monthly climatology calculated over 1999-2009, on a $\frac{1}{4}^\circ$ grid
 134 [Lungu and Callahan, 2006]. ROMS is connected to the lateral boundaries by an active, implicit
 135 and upstream-biased radiation condition [Marchesiello et al., 2001]. The boundary conditions
 136 and initial state are based on the objectively analysed World Ocean Atlas 2013 (WOA13)
 137 monthly climatology on a $\frac{1}{4}^\circ$ grid [Locarnini et al., 2013; Zweng et al., 2013]. The inflow
 138 boundary conditions are nudged toward temperature, salinity, and geostrophic velocity fields.
 139 Following Kersalé et al., [2011], the nudging timescale for inflow and outflow are set to 1 day
 140 and 1 yr for the tracer fields and 3 days and 1 yr for the momentum fields. The geostrophic
 141 velocity is referenced to the 1000 m depth. The explicit lateral viscosity is null all over the
 142 domain, except in the sponge layer. The width of the nudging border is 50 km and the maximum
 143 viscosity value for the sponge layer is set to $1000 \text{ m}^2 \text{ s}^{-1}$.
 144 The high resolution ROMS simulation, hereafter referred to as W13Q, is run over 10 years with
 145 outputs averaged every 2 days. Integrated physical properties show that a statistical equilibrium
 146 is reached by the model after the first year of simulation. To avoid any impact of the spin up on
 147 the model output and to ensure the robustness of our findings, we decided to also entirely remove
 148 the third year and to focus our investigation on the outputs from year 4 to 10.

149 2.2 Datasets

150 We compared our numerical results with several datasets. The Moderate Resolution Imaging
 151 Spectroradiometer (MODIS) infrared SST measured from the Aqua satellite are available on a
 152 monthly basis with a spatial resolution of 4 km. We use the time period from July 2002 to June
 153 2015 to compute the monthly climatology.

154 The 3-D monthly fields of temperature and salinity issued from the In Situ Analysis System
 155 (ISAS13) are based on Argo data [Gaillard et al., 2016] and are available on a 0.5° grid from
 156 2002 to 2012.

157 Monthly near surface current over 1992 to 2015 are obtained from the Ocean Surface Current
 158 Analysis – Real time (OSCAR) with a $1/3^\circ$ spatial resolution [Bonjean and Lagerloef, 2002].

159 Finally we used reanalyzed currents from the global Hybrid Coordinate Ocean model and the
 160 Navy Coupled Ocean Data Assimilation (HYCOM+NCODA) [Cummings and Smedstad, 2013].
 161 These velocity fields are available with a spatial resolution of $1/12.5^\circ$ and a daily resolution from
 162 January 1st 2006 to December 31st 2012.

163 2.3 Non-dimensional numbers

164 The ocean dynamics in island wake is generally characterized by the recourse to non-
 165 dimensional numbers issued from geophysical fluid dynamics. The wake is typically controlled
 166 by the turbulent Reynolds number defined as:

$$Re = U_0 D / \nu$$

167 where U_0 is the unperturbed upstream velocity, D is the horizontal scale of the obstacle, and ν is
 168 the eddy viscosity [Tomczak, 1988]. To compute Re in the Marquesas Islands, we use an
 169 horizontal eddy viscosity value of $\nu = 100 \text{ m}^2 \text{ s}^{-1}$ [Heywood et al., 1990; Jiménez et al., 2008].
 170 In the archipelago, only the biggest island, Nuku Hiva, exhibits a Reynolds number ($Re \approx 50$)
 171 that exceeds the theoretical threshold of the Von Karman vortex street generation ($Re_{th} = 40$).

172 To consider the spatial scale and energy involved in island wake, we also consider the Burger
 173 number, representing the ratio between stratification and Earth's rotation, and defined as:

$$Bu = \left(\frac{R_d}{R}\right)^2$$

174 where R is the island radius and R_d the Rossby radius associated to the first baroclinic mode.
 175 Using the forcing data from WOA13, we found that the seasonality of R_d vary between 120 and
 176 130 km in the archipelago. When the island radius R is smaller than R_d ($Bu \geq 1$), a submesoscale
 177 wake is generated with eddies having their radius in the order of R [Stegner, 2014]. This is
 178 always the case in the archipelago where the radius of the largest island is about 20 km.

179 Finally, we calculate the geometric shallow-water parameter $\alpha = \frac{h}{R^2}$, with h the mean thermocline
 180 depth. We obtained an alpha ranging between 0.01 and 0.03. These values are typical of mid-
 181 ocean isolated islands, such as Madeira [Caldeira *et al.*, 2002], Gran Canaria [Sangrà *et al.*,
 182 2005] or Hawaii [Calil *et al.*, 2008], generating a deep-water wake.

183 2.4 Eddy detection and tracking

184 To detect eddies in our model experiment, the method based on the geometry of the flow field
 185 developed by Nencioli *et al.* [2010] is applied. This method is well established and used in
 186 several studies with different resolution [Dong *et al.*, 2012; Liu *et al.*, 2012; Amores *et al.*, 2013;
 187 Mkhinini, N. *et al.*, 2014]. It relies on the spatial definition of vortices defined as a region with a
 188 rotary flow around its center. In other words, a vortex is defined as a region with velocity vectors
 189 rotating around a center with a minimum speed. Its boundary is defined by the largest closed
 190 contour line of the local stream function around the center. The flexibility of the algorithm
 191 depends on two parameters: a , the number of grid points where the current increases, and b , the
 192 dimension of the area used to define the local minimum of velocity. These parameters set the
 193 minimum size of the detectable vortices and allow the algorithm to work on different grid
 194 resolution. Using the same protocol than Nencioli *et al.* [2010] and Liu *et al.* [2012], optimal
 195 performances of the algorithm have been obtained with $a = 2$; $b = 2$ for our model
 196 configuration and $a = 3$; $b = 2$ for the HYCOM reanalysis outputs.

197 Once eddies have been detected at each time step, we used the tracking method proposed by
 198 Doglioli *et al.* [2007] and Nencioli *et al.* [2010]: the position and sign of each eddy center are
 199 compared at successive time steps. An eddy track is identified when the eddy center at the next
 200 time step is found within a given searching area of the previous time step. The searching radius is
 201 the product of the averaged eddy center displacement speed and the time interval between two
 202 model outputs (i.e., 2 days). To avoid that eddies move further than the searching area within
 203 successive time steps (thus preventing track splitting), we added a constant of 20 cm s^{-1}
 204 corresponding to the eddy-center displacement speed and representing the mean velocity of the
 205 background current. Moreover, eddies with a shorter lifetime than 4 days are not considered to
 206 ensure the only detection of consistent structures only.

207 The dataset obtained after applying the detection includes: eddy center locations, eddy contours,
 208 polarities and radii. Time evolution of each eddy detected is also recorded and gives us
 209 information on their generation and ending.

210 2.5 Eddy Kinetic Energy budgets

211 The *EKE* is defined as the kinetic energy due to transient dynamic. Neglecting the vertical
 212 velocity contributions: $EKE = \frac{1}{2}(u'^2 + v'^2)$, where $u' = u - \bar{u}$ and $v' = v - \bar{v}$, with u and v
 213 the zonal and meridian components of the velocity at 10 m, \bar{u} and \bar{v} their temporal means from
 214 year 4 to 10. To consistently compare our numerical model results with the monthly satellite

215 derived observations (OSCAR), we subsampled the W13Q current components at 10 m (u_m and
 216 v_m) by extracting one grid point every 5 grid points. Then, we used monthly averages to
 217 compute the monthly $EKE_m = \frac{1}{2}(u_m'^2 + v_m'^2)$.

218 To quantify the relative importance of instability and eddy-mean interaction mechanisms on the
 219 EKE generation, we consider the different terms of energy budgets relative to the EKE
 220 generation [Auad et al., 1991; Marchesiello et al., 2003; Liang et al., 2012]. The wind work, the
 221 barotropic and the baroclinic energy conversion [Dong et al., 2007; Halo et al., 2014; Hristova et
 222 al., 2014; Sun et al., 2016] are defined, respectively, as:

$$\begin{aligned}
 FeKe &= -\frac{1}{\rho_0}(\overline{u'\tau_x'} + \overline{v'\tau_y'}) \\
 KmKe &= -\left(\overline{u'u'}\frac{\partial\bar{u}}{\partial x} + \overline{u'v'}\frac{\partial\bar{u}}{\partial y} + \overline{u'w'}\frac{\partial\bar{u}}{\partial z} + \overline{v'u'}\frac{\partial\bar{v}}{\partial x} + \overline{v'v'}\frac{\partial\bar{v}}{\partial y} + \overline{v'w'}\frac{\partial\bar{v}}{\partial z}\right) \\
 PeKe &= -\frac{g}{\rho_0}\overline{\rho'w'}
 \end{aligned}$$

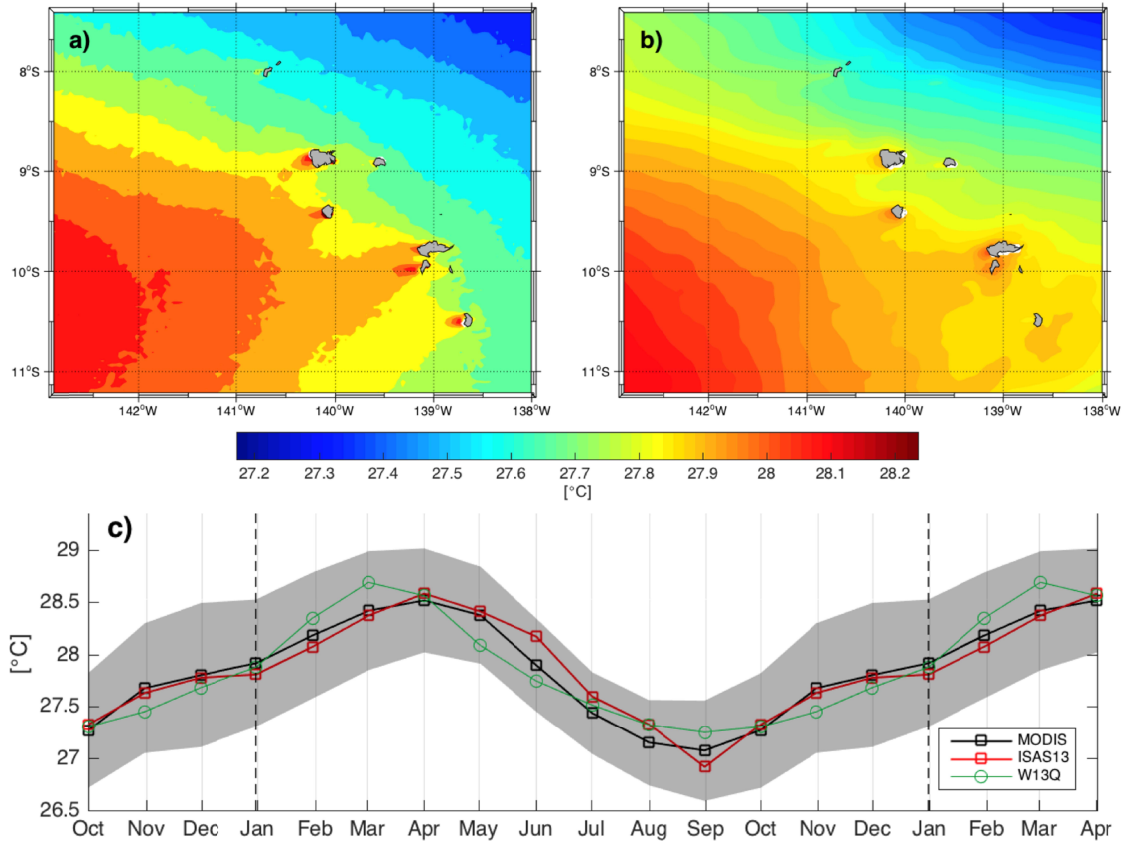
223 where τ_x and τ_y are the meridian and zonal component of the wind stress. The prime indicates
 224 the anomaly from the annual mean. g is the gravity, ρ is the density and $\rho_0 = 1030 \text{ kg m}^{-3}$. The
 225 overbar denotes the temporal mean and the prime represents the anomaly deviation from the
 226 mean. The $FeKe$ is the EKE generation due to transient wind. When positive, this term induces
 227 an energy input while, when negative, it implies a damping effect on the sea surface [Xu and
 228 Scott, 2008]. The $KmKe$ is the energy conversion between mean currents and EKE . When its
 229 volume-integration is positive, it implies barotropic energy conversion. The $PeKe$ is the energy
 230 conversion between available potential energy and EKE and indicates baroclinic instability when
 231 its volume-integration is positive [Harrison and Robinson, 1978]. Both $KmKe$ and $PeKe$ are
 232 vertically integrated over the surface layer (top 100 m) [Hristova et al., 2014]. The minimum
 233 depth in the model is set to 75 m. When the bathymetry is shallower than 100 m, conversion
 234 terms are integrated over the actual depth.
 235

236 3 Results and Discussion

237 3.1 Thermohaline structures and seasonal variability

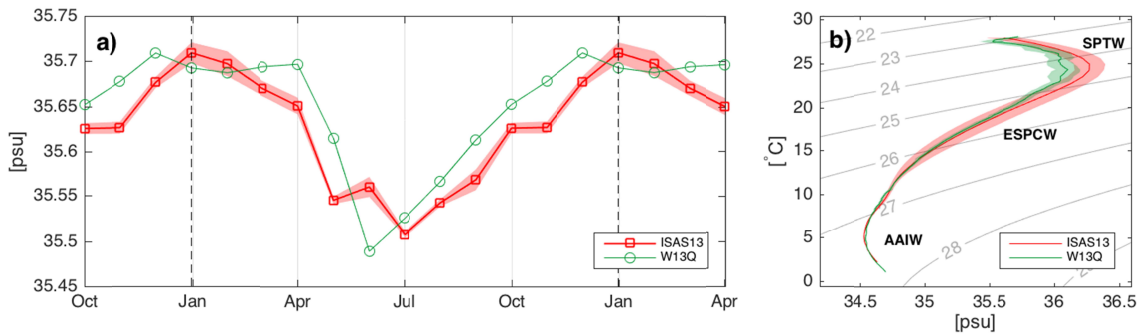
238 The Marquesas archipelago is characterized by a north-east/south-west SST gradient as shown
 239 by MODIS remote sensing data and W13Q (Figure 3a and 3b). Colder water in the north-east (\leq
 240 27.5°C) represents the imprint of the equatorial upwelling [Wyrski, 1981] while warmer SST in
 241 the west ($\geq 28^\circ\text{C}$) is the south-eastward imprint of the western Pacific warm pool [Yan et al.,
 242 1992]. The high-resolution of MODIS reveals warm regions behind the islands likely due to
 243 areas zones of weak currents leeward the islands. These structures have also been observed
 244 leeward several islands elsewhere [Caldeira and Marchesiello, 2002; Caldeira et al., 2002]. The
 245 maximum of the mean satellite SST over the archipelago is up to 28.6°C in April during the
 246 austral summer (Figure 3c, black line). In austral winter (September), SST in the Marquesas
 247 region decreases down to 27°C . In this tropical region, the seasonal amplitude is around 2°C as
 248 confirmed by *in-situ* SST from ISAS13 (red line). All the aforementioned features appear in the
 249 same locations and with comparable amplitudes in our numerical experiment, although SST is
 250 slightly colder than satellite observations (0.05°C to 0.1°C ; Figure 3b vs. 3a; green line on Figure

251 3c). Thus, SST seasonal climatology from the model shows a good agreement with *in situ* and
 252 satellite-derived temperature variability.



253 Figure 3: Mean SST over the child grid area from (a) MODIS (time averaged from Jul-2002 to
 254 Jun-2015) and (b) W13Q at 10m (time averaged over year 4 to 10). (c) Monthly climatology
 255 averaged over the child grid area from MODIS (black), ISAS13 (red) and W13Q (green). The
 256 grey shaded area represents the spatial standard deviation from MODIS monthly climatology.
 257 Units are in °C.
 258
 259

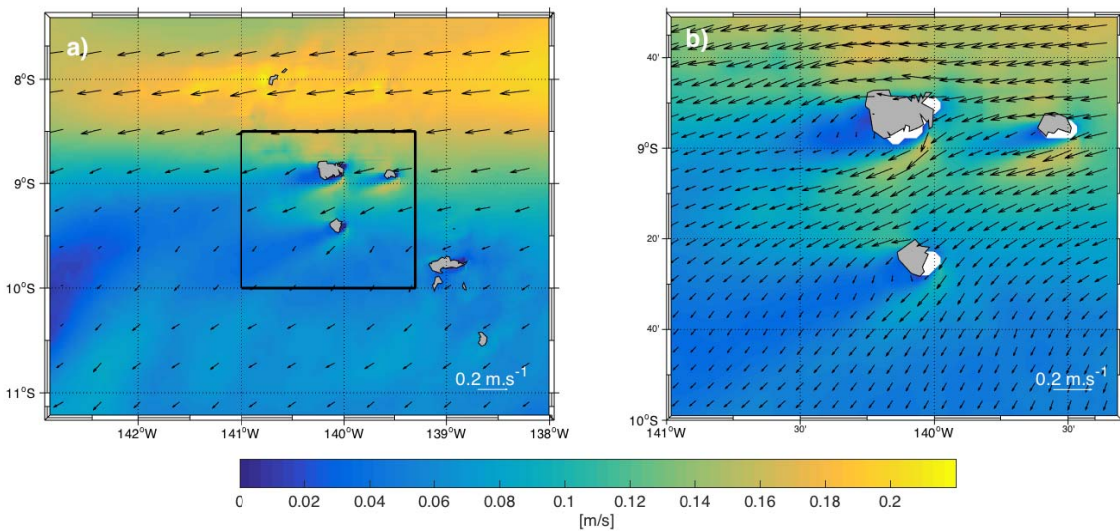
260 Considering the haline structure, *in situ* data from ISAS13 show a SSS maximum in austral
 261 summer and a minimum in austral winter (35.7 and 35.5 psu +/- 0.01 psu), that are satisfactory
 262 reproduced by the model (Figure 4a). The TS diagram issued from ISAS13 and W13Q illustrates
 263 the water masses within the archipelago (Figure 4b). The South Pacific Tropical Water (SPTW)
 264 is characterized by a salinity maximum in subsurface (> 35.6 psu) [O'Connor et al., 2005]. The
 265 Eastern South Pacific Central Water (ESPCW) and the Antarctic Intermediate Water (AAIW) are
 266 respectively defined by temperature ranges between 8 °C and 24 °C and between 2 °C and 10 °C.
 267 Their salinities are ranging between 34.4 psu and 36.4 psu and between 33.8 psu and 34.5 psu,
 268 respectively [Emery, 2001]. Only the SPTW is fresher in W13Q than with *in situ* measurements.
 269 The other water masses are quite well reproduced.



270
 271 Figure 4: (a) Time series of the monthly averaged salinity (in psu) at 10 m issued from ISAS13
 272 (red) and W13Q (green). Red shaded area corresponds to the spatial standard deviation from
 273 ISAS13. (b) TS diagram from ISAS13 (red) and W13Q (green) monthly climatology and their
 274 standard deviations (red and green shades respectively).
 275

276 **3.2 Current structures**

277 As shown in Figure 1a, the SEC is the main surface current flowing around the Marquesas
 278 archipelago. Currents issued from W13Q present the same order of magnitude than the OSCAR
 279 ones, although with a more westward direction in the north (Figure 5a). The spatial distribution
 280 of the mean current in the archipelago can be depicted in two regions: the southern part where
 281 the horizontal current is weak (8 cm s^{-1}), and the northern part where the SEC is stronger and
 282 reaches values of 20 cm s^{-1} as observed by *Martinez et al.* [2009]. The deviation of the SEC
 283 creates stronger currents on both sides of the islands and regions of weak current just behind the
 284 islands (Figure 5b) [*Chang et al.*, 2013; *Karnauskas et al.*, 2017]. These shadow zones of weak
 285 currents are collocated with the warm areas as reported previously in *Figure 3b*.
 286

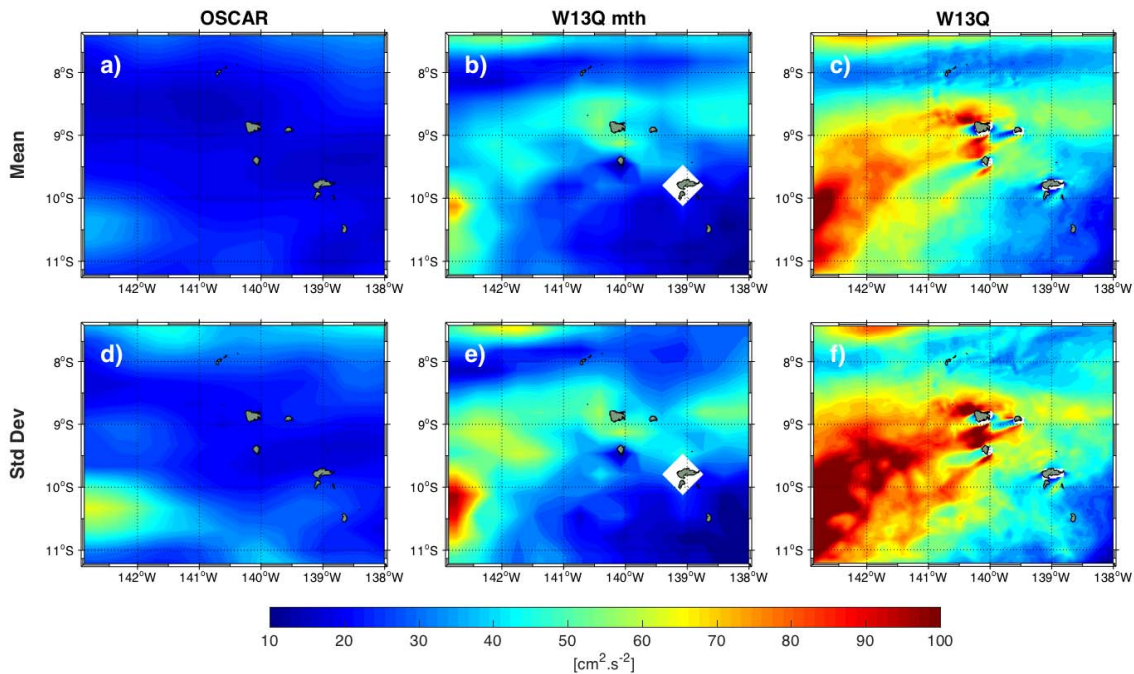


287
 288 Figure 5: (a) W13Q current at 10 m time averaged over year 4 to 10 and over the child domain.
 289 For clarity, only 1 vector out of 17 is represented. (b) Zoom of (a) over the Marquesas northern
 290 islands, as defined by the black box on a). For clarity, only 1 vector out of 3 is represented. Units
 291 are in m s^{-1} .

292

293 3.3 *EKE* generation and seasonal variability

294 In order to investigate the dynamical properties within the archipelago and more precisely the
 295 eddy field, we first focus on the *EKE* characterisation. The EKE_m obtained from OSCAR
 296 monthly currents show homogeneously low values around the archipelago ($25 \text{ cm}^2 \text{ s}^{-2}$) in Figure
 297 6a. In the northern and the south-western part of the archipelago, two areas are barely more
 298 active, the former corresponding to the southward imprint of the turbulent equatorial area [Qiu
 299 and Chen, 2004]. To compare the eddy activity from W13Q with the remote sensed one, we
 300 subsampled the model outputs to the same spatial and temporal resolution before computing
 301 EKE_m (Figure 6b). High activity occurs in the northern and south-western part of the
 302 archipelago. Monthly W13Q also presents a high eddy activity west of Nuku Hiva. To fully take
 303 advantage of the high-resolution model, the annual mean of *EKE* is now directly computed from
 304 the W13Q two-day outputs (Figure 6c). Thanks to the higher spatial and temporal resolution, the
 305 *EKE* at 10 m reaches up to $90 \text{ cm}^2 \text{ s}^{-2}$ (vs. $55 \text{ cm}^2 \text{ s}^{-2}$ in Figure 6b). Island wakes characterized by
 306 a highly variable *EKE* pattern are highlighted westward the northern islands of Nuku Hiva, Ua
 307 Pou and Ua Huka. We could have also expected high activity leeward Hiva Oa (in the south),
 308 since its diameter is approximately the same as Nuku Hiva. However, the SEC being weaker
 309 windward this island, it does not produce a strong enough *EKE*. Consistently, *EKE* patterns are
 310 more pronounced in the northern part of the archipelago where the mean current is stronger as
 311 seen on Figure 5a. The *EKE* standard deviations are given in Figure 6e and 6f. They reveal high
 312 activity where the mean *EKE* is already high and in good agreement with OSCAR *EKE* standard
 313 deviations (Figure 6d).



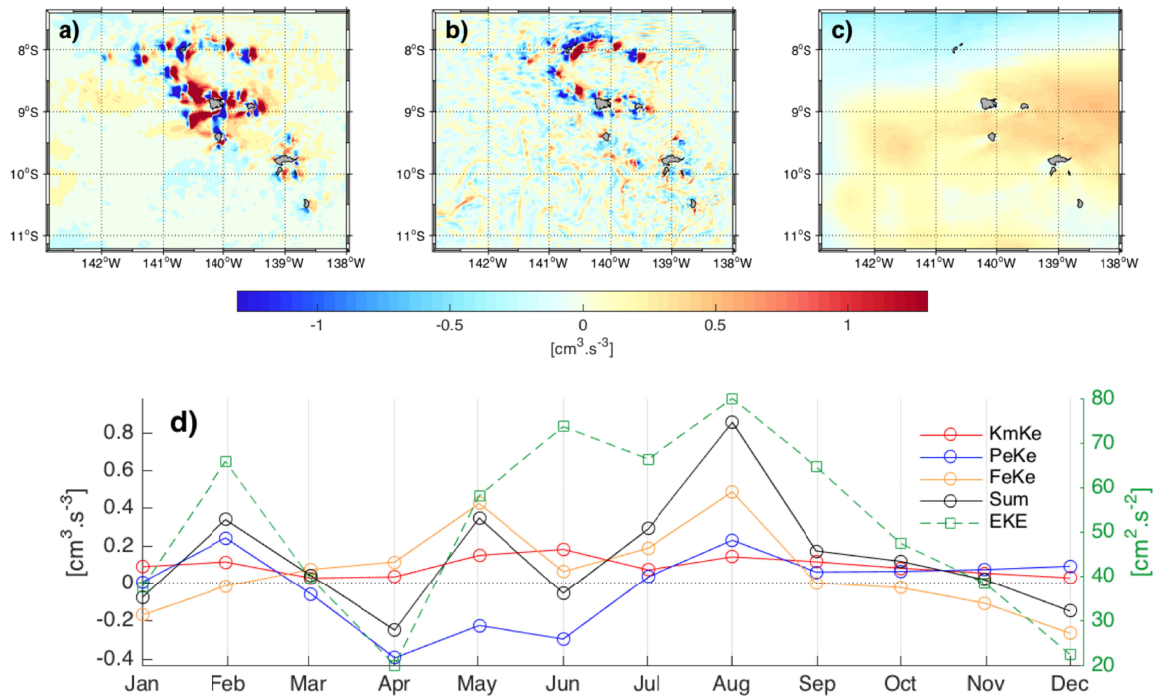
314

315 Figure 6: Annual mean of EKE_m (in $\text{cm}^2 \text{ s}^{-2}$) obtained from monthly averaged of (a) OSCAR
 316 surface currents and (b) W13Q currents at 10m. (c) is the annual mean of *EKE* derived from the
 317 W13Q 2-day outputs. Their standard deviations are represented in (d), (e) and (f), respectively.

318

319 In order to determinate the *EKE* origin, we now investigate the energy budget. For all the islands,
320 the Burger number $Bu > 1$. Therefore, a significant amount of kinetic energy in the *EKE* origin is
321 expected. This trend is observed in the barotropic energy conversion (Figure 7a). The *KmKe* is
322 particularly intense leeward Nuku Hiva where the *EKE* pattern is the strongest (Figure 6c). It is
323 mainly due to the fact that Nuku Hiva is the largest island and that it lies where the currents are
324 the strongest. This implies a $Re = 58$ and, consequently, a possible Von Karman wake
325 generation. *KmKe* and *PeKe* spatial patterns are close and are driven by the topography of the
326 archipelago (Figure 7a, b and Figure 1b). The distribution of the baroclinic energy conversion is
327 negative in average, which indicates an eddy dissipation toward the eddy potential energy [*Kang*
328 *and Curchitser, 2015*]. The averaged energy input from the wind (*FeKe*) presents smooth
329 patterns because of the relatively low resolution of the wind stress used in the model (Figure 7c).
330 Values are weaker than the two previous parameters.

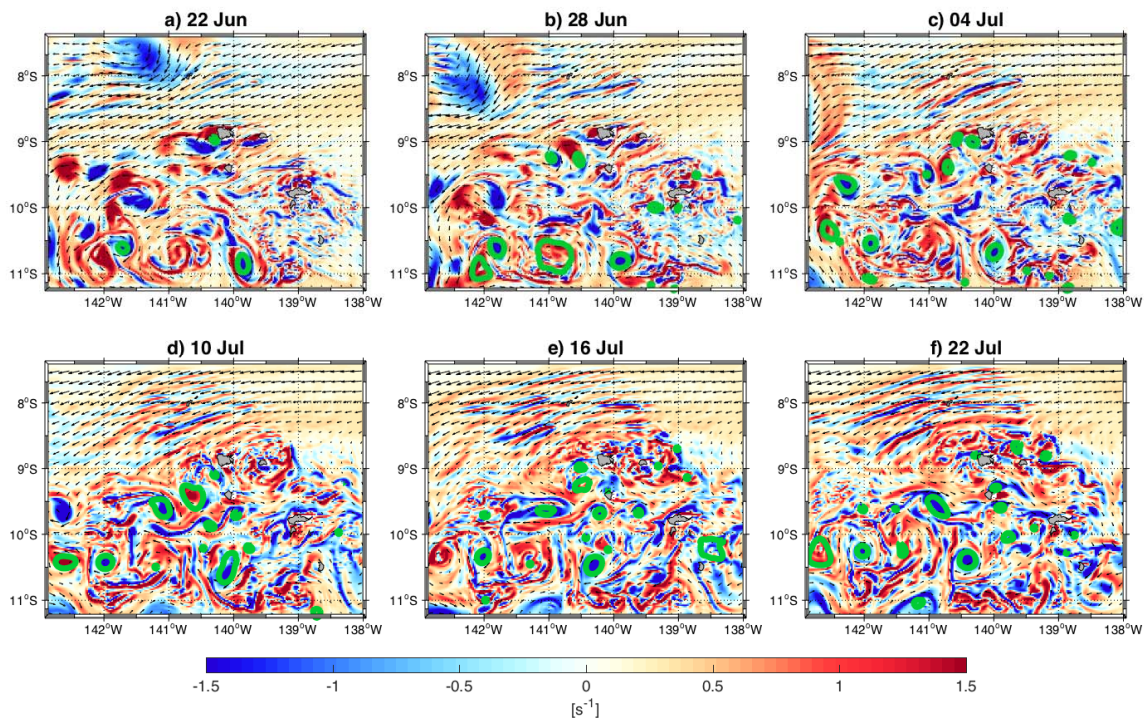
331 The *EKE* seasonality is superimposed with the different energy budget terms and their sum in
332 Figure 7d. *EKE* time series show a maximum activity in February, June and August, when the
333 *KmKe* is also high. The latter remains positive all year round, indicating a constant flux towards
334 *EKE*. *PeKe*, in a smaller proportion, also contributes to the high *EKE* activity in February. *PeKe*
335 is responsible for the minimum value of *EKE* in April. The *EKE* variation on the second half of
336 the year is explained by the *FeKe* variation, with a decreasing activity from August to December.
337 The sum of these energy flux terms is relatively well correlated with the *EKE* variability. When
338 this sum is positive (negative), an increase (decrease) of *EKE* is observed except during June
339 where the energy budget seems too low to explain the *EKE* activity. According to *Chen et al.*
340 [2016], the energy transfer through eddy-mean flow interactions could be imbalanced and
341 compensated by advection. Indeed, currents are relatively intense in the archipelago during this
342 time period. This could be at the origin of the decay observed in these time series. Although the
343 spatial averaging reveals low values of *PeKe* and *FeKe* compared to some local maxima (Figure
344 7a and b), these energy inputs seem to play an important role in the *EKE* seasonal variation,
345 particularly at the beginning and the end of the year.
346



347 Figure 7: Spatial distribution of the annual mean of the (a) barotropic ($KmKe$), (b) baroclinic
 348 ($PeKe$) energy conversion both integrated over the surface to 100 m and (c) the EKE generation
 349 due to transient wind, from year 4 to 10 of W13Q (units are in $cm^3 \cdot s^{-3}$). (d) Time series of wind
 350 work, baroclinic and barotropic energy (left axis), in yellow, blue and red, respectively. The sum
 351 of these three terms is indicated in black. The EKE (in $cm^2 \cdot s^{-2}$, right axis) spatially averaged over
 352 the child grid and at 10 m is represented in green.
 353

354 3.4 Eddy activity

355
 356 Detected eddies within the Marquesas archipelago are illustrated with green contours on relative
 357 vorticity maps from June 22 to July 22 during year 9 of W13Q (Figure 8). Positive (negative)
 358 values reveal anticyclonic (cyclonic) circulation. The total number of detected eddies over the 7
 359 years of integration is 1260, with a dominance of cyclonic (714) vs. anticyclonic eddies (546).
 360 This cyclonic eddy dominance (i.e., 56.7% of the total eddy number) is well illustrated (Figure
 361 8). On the other hand, positive patterns are unstable and tend to broaden and fragment inducing a
 362 weaker proportion of anticyclonic eddies. These observations are consistent with previous
 363 studies and are due to centrifugal instability occurring to anticyclonic eddies [*Dong et al.*, 2007;
 364 *Hasegawa et al.*, 2009; *Stegner*, 2014]. These snapshots reveal an important eddy activity in the
 365 archipelago which has never been investigated before. In particular, Figure 8c shows the
 366 generation of cyclonic and anticyclonic eddies leeward Nuku Hiva, the northern island, and
 367 likely producing SST small-scale patterns as observed in Figure 2.

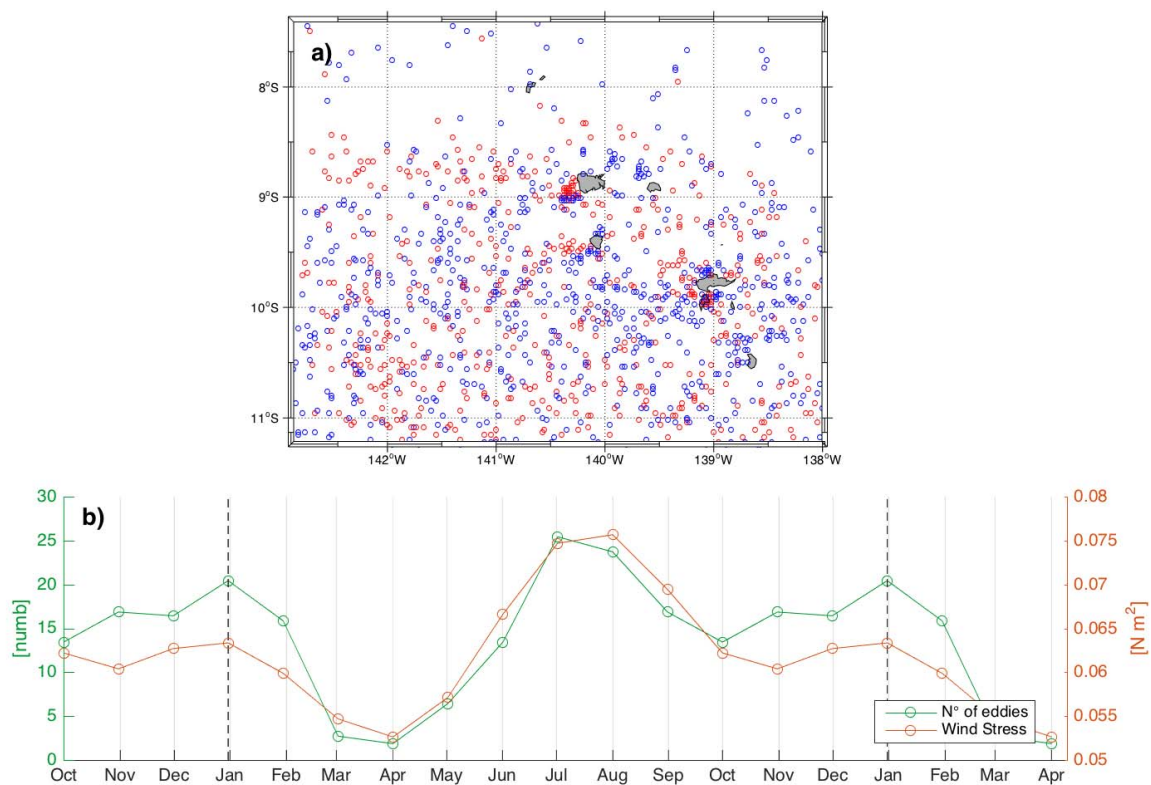


368
 369 Figure 8: Snapshots of relative vorticity at 10 m (units are in s^{-1}) from W13Q during Year 9
 370 every 6 days (a-f). The detected eddies are represented by *green* contours (Nencioli et al. [2010]
 371 based algorithm). The animation of the vorticity from June to July is available in the auxiliary
 372 material.

373
 374 To investigate where eddies are generated, the first eddy position determined by the detection
 375 algorithm is recorded and analysed. Most of eddies are generated leeward the archipelago
 376 highlighting the role of the topography (Figure 9a). Part of these eddies are generated just behind
 377 the islands. Although the mean Reynolds number is larger than 40 only for Nuku Hiva, eddies
 378 are also generated behind the other islands. Indeed, when surface currents accelerate on the
 379 flanks of the other islands (Figure 5), they induce a larger Re and create a decreasing pressure in
 380 the lee side associated with two opposite recirculation cells behind the islands. Anticyclonic
 381 eddies are generated on the equatorward side and cyclonic eddies on the poleward side. These
 382 eddies are emphasized behind Nuku Hiva and Ua Pou in the northern part of the archipelago
 383 where the SEC is the strongest, and Hiva Oa in the south which is the second biggest island of
 384 the Marquesas archipelago. Conversely, there is no eddy generation behind Ua Huka (the
 385 northeastern most island) due to its small size inducing a too low Re ($= 34$) while there is a
 386 signature of the island wake on EKE (Figure 6c), mean currents (Figure 5), energy conversion
 387 rates (Figure 7a) and vorticity (Figure 8).

388 Eddies are also generated farther from the coasts, in the open ocean, where the flow presents
 389 complex structures as seen on Figure 8. They are generated almost exclusively south-westward
 390 of the islands. A small number is also generated elsewhere, likely due to an eastward
 391 perturbation of the SEC during austral winter and referred as the Marquesas Counter Current
 392 [Martinez et al., 2009]. Most of the eddies in the archipelago are generated during austral winter
 393 (July and August) when the wind stress and the EKE are maximum (Figure 9b and Figure 7c).
 394 Contrarily, eddies are barely detected in March and April when the wind stress and EKE are

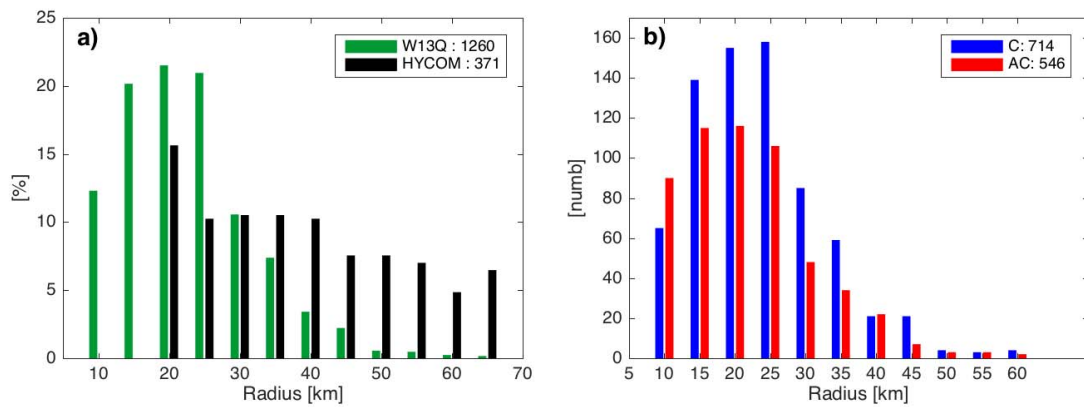
395 minimum. A strong correlation is found between eddy generation and the wind stress, with
 396 $r = 0.87$ ($p = 2.10^{-4}$). A high eddy generation occurs when the wind is strong (Figure 9b).
 397 Indeed, not only the wind is responsible for a part of the *EKE* generation, but it also strengthens
 398 the surface currents, providing sufficient velocity to generate eddies behind the islands.



399 Figure 9: (a) Location of eddy generation identified from W13Q over year 4 to 10. Anticyclonic
 400 and cyclonic eddies are represented by red and blue circles, respectively. (b) Monthly mean
 402 number of generated eddies from W13Q from year 4 to 10 (green line) and wind stress issued
 403 from QuickSCAT in $N \cdot m^{-2}$ (orange line).
 404

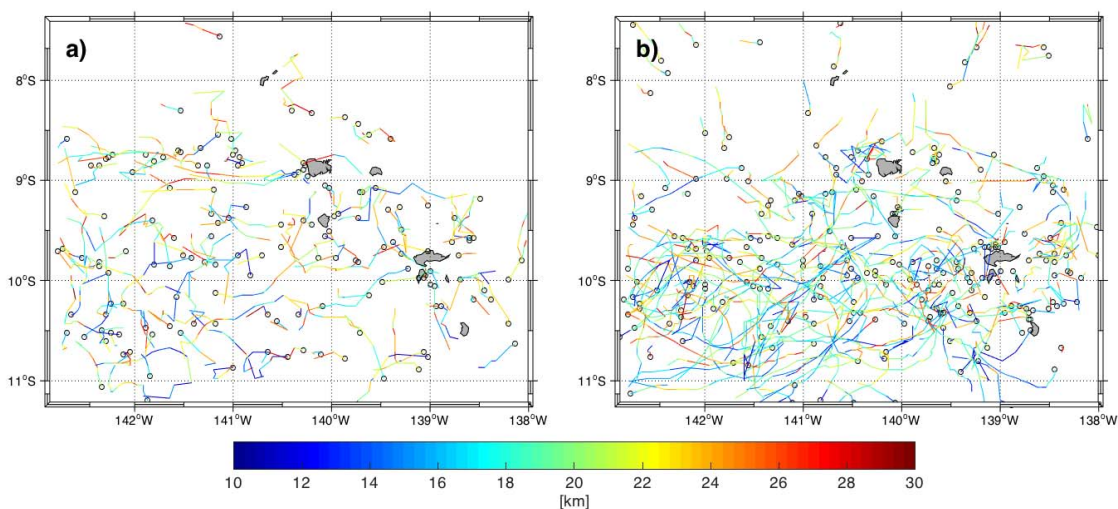
405 Investigating the maximum size of the detected eddies during their lifetime in W13Q shows that
 406 their distribution is similar to a Gaussian without an inferior tail due to resolution limitations
 407 (Figure 10a, green bars). Most of the eddy radii are lower than 25 km and only a few are larger
 408 than 45 km. Indeed, size of eddies generated leeward the islands are approximately the same than
 409 the island dimensions, which can be explained by Bu larger than 1. The HYCOM eddy size
 410 distribution from January 2006 to December 2012 has been investigated within the child grid
 411 (Figure 10a, black bars). In opposition with the W13Q results, a significant number of eddies
 412 larger than 20 km is present while smaller eddies are absent in HYCOM. The coarser spatial
 413 resolution of HYCOM prevents the resolution of such small structures. HYCOM larger eddies
 414 are suspected to be induced by the inter-annual forcing which is not represented in the present
 415 climatological W13Q configuration. In total, almost three times more eddies, are detected in our
 416 configuration than in HYCOM resulting from the higher resolution of our model.
 417 Considering now the cyclonic vs. anticyclonic patterns, W13Q cyclonic eddies are dominant for
 418 all radius classes, except for the smallest ones (~ 10 km) (Figure 10b). Nevertheless, as previously

419 mentioned, the generated anticyclonic structures are more unstable and form small structures
 420 with short lifetime (shorter than a week).



421
 422 Figure 10: (a) Distribution of the detected eddy radius from year 4 to 10 in W13Q (green) and
 423 from 2006 to 2012 in HYCOM (black). (b) Distribution of W13Q cyclonic (blue) vs.
 424 anticyclonic (red) eddies.
 425

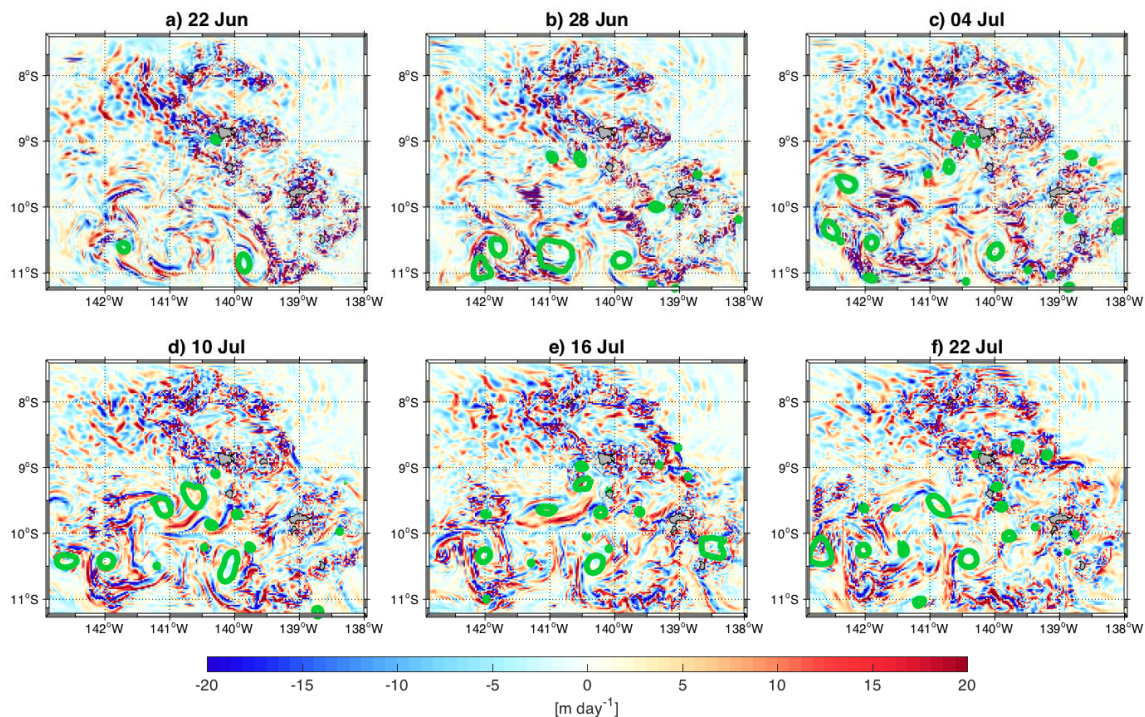
426 Focusing on these small eddies of interest, we further investigate trajectories of eddies with a
 427 maximum radius between 20 and 30 km, which represents about 50% of the eddies (Figure 10a).
 428 As presented in Figure 10b, anticyclonic eddies are less numerous than cyclonic ones for this
 429 radius range (Figure 11a vs. 11b, respectively). Furthermore, the anticyclonic eddy trajectories
 430 are shorter than the cyclonic ones (Figure 11b) because of the centrifugal instability that weakens
 431 anticyclonic eddies [Dong *et al.*, 2007]. While Calil *et al.* [2008] showed that cyclonic
 432 (anticyclonic) eddies tend to propagate poleward (equatorward) in the Hawaiian archipelago,
 433 trajectories of eddies generated leeward the Marquesas Islands are mostly zonal and do not
 434 present such a meridian trend. Indeed, while 61% (54.8%) of the cyclonic (anticyclonic) eddies
 435 propagate equatorward (southward), their meridian mean latitude change is only of 0.0163°
 436 (0.0182°). This might be explained by the lower latitude of the Marquesas archipelago inducing a
 437 weaker beta effect than in Hawaii.



439 Figure 11: Eddy generation sites (black circles) and their trajectories for (a) anticyclonic and (b)
 440 cyclonic eddies detected in W13Q and for a maximum radius over 20-30 km encountered over
 441 their lifetime. The size of the radius along the trajectory is given by the color bar (in km).
 442

443 3.5 An insight to a possible eddy induced biological activity

444 *Hasegawa et al.* [2009] hypothesized that a combination of three processes can enrich the
 445 surface layer in island wakes: the propagation of coastal rich waters formed in the lee of the
 446 islands by eddy shedding, the vertical advection of the Deep Chlorophyll Maximum (DCM)
 447 toward the surface in the island wake, and the uplift of nutrient rich deep waters by the strongest
 448 cyclonic eddies. We show that, when the SEC encounters the Marquesas Islands, it generates
 449 eddies in the lee of the archipelago (Figure 9a). These eddies can trap and propagate inshore
 450 waters in the island wake. Thus, if these inshore waters have been previously enriched in
 451 proximity of the islands, the biological activity can be enhanced. An increase of the sea surface
 452 concentration of chlorophyll in the lee of an island could also be induced by the vertical
 453 advection of the DCM. During the “Pakaihi i te Moana” expedition, a correlation between the
 454 depth of the DCM and the mixed layer depth (MLD) has been reported in the archipelago
 455 [Martinez et al., 2016]. On average, the DCM is deeper (shallower) in the northern (southern)
 456 region of the archipelago consistently with the MLD (60 m vs. 30 m). Our results show strong
 457 vertical velocities at the MLD associated with the eddies of the island wake (Figure 12). This
 458 high vertical activity might induce the vertical advection of the DCM toward the surface. The
 459 evaluation of the impact of the fine scale dynamics on the nutricline is beyond the scope of the
 460 present study and will need further research. Indeed, to assess the uplift of nutrient-rich deep
 461 waters by cyclonic eddies and the associated biological activity, a coupled physical-
 462 biogeochemical model is required.
 463



465 Figure 12: Snapshots of vertical velocities (in m day^{-1}) as in Figure 8, but at the mixed layer
466 depth.
467

468 **4 Conclusions and perspectives**

469 The island wake in the Marquesas archipelago has been characterized in the present study,
470 providing new insights into the spatial and seasonal variability of the fine-scale dynamics of this
471 oceanic region. We based our analysis on the ocean circulation model ROMS-AGRIF with a
472 $1/45^\circ$ child grid resolution and forced by climatological mean fields. The high spatial and
473 temporal resolution allows characterizing warm SST patterns just behind the Marquesas Islands.
474 These patterns likely due to shadow zones of weak currents have been consistently observed
475 leeward several islands [Barton, 2001; Caldeira et al., 2002; Chavanne et al., 2002].

476 Accordingly to a rough approximation of the local Reynolds number, only Nuku Hiva, the
477 biggest island of the archipelago, would be expected to generate a wake of propagating eddies.
478 However, the mean *EKE* and the use of an automated eddy detection algorithm reveals that eddy
479 generation also occurs leeward the other islands, inshore and offshore. The energy budget
480 analysis reveals that eddies are generated by a combination of wind power input, barotropic and
481 baroclinic energy conversion. These generated eddies are cyclonic dominant due to the small size
482 of the islands in comparison to the Rossby radius. Most of the eddies have their maximum radius
483 between 20 and 30 km.

484 Our physical model results provide insight to explain the Marquesas biological enhancement. We
485 suggest that the combination of at least two of the three processes proposed by Hasegawa et al.
486 [2009] enriches surface layers via eddy activity: the propagation of rich waters formed in the lee
487 of the islands by eddy shedding and the vertical advection of the DCM toward the surface in the
488 island wake. To further investigate the third process, I.E. the possibility of an uplift of nutrient
489 rich waters by the cyclonic eddies, new in situ data and the use of a coupled physical-
490 biogeochemical numerical model are required.
491

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496 The data used in this work are available from Moderate Resolution Imaging Spectroradiometer
497 MODIS (<http://oceandata.sci.gsfc.nasa.gov>), In Situ Analysis System ISAS13
498 (<http://www.seanoe.org/data/00348/45945/>), Ocean Surface Current Analysis – Real time
499 OSCAR (<http://www.oscar.noaa.gov/>) and Hybrid Coordinate Ocean Model HYCOM
500 (<http://hycom.org/data/glb0pt08/expt-19pt1>). Please contact the corresponding author to obtain
501 the data of the numerical model used in this paper.
502

503 **References**

- 504 Alae, M. J., G. Ivey, and C. Pattiaratchi (2004), Secondary circulation induced by flow
505 curvature and Coriolis effects around headlands and islands, *Ocean Dyn.*, 54(1), 27–38,
506 doi:10.1007/s10236-003-0058-3.
507 Amores, A., S. Monserrat, and M. Marcos (2013), Vertical structure and temporal evolution of

- 508 an anticyclonic eddy in the Balearic Sea (western Mediterranean), *J. Geophys. Res. Ocean.*,
509 118(4), 2097–2106, doi:10.1002/jgrc.20150.
- 510 Andrade, I., P. Sangrà, S. Hormazabal, and M. Correa-Ramirez (2014), Island mass effect in the
511 Juan Fernández Archipelago (33°S), Southeastern Pacific, *Deep. Res. Part I Oceanogr. Res.*
512 *Pap.*, 84, 86–99, doi:10.1016/j.dsr.2013.10.009.
- 513 Auad, G., A. Pares-Sierra, and G. K. Vallis (1991), Circulation and energetics of a model of the
514 California Current system.pdf, *J. Phys. Oceanogr.*, 21(10), 1534–1552.
- 515 Barton, E. D. (2001), Island wakes, *Nature*, 1–8, doi:10.1006/rwos.2001.0140.
- 516 Beckmann, A., and D. B. Haidvogel (1993), Numerical Simulation of Flow around a Tall
517 Isolated Seamount. Part I: Problem Formulation and Model Accuracy, *J. Phys. Oceanogr.*,
518 23(8), 1736–1753, doi:10.1175/1520-0485(1993)023<1736:NSOFAA>2.0.CO;2.
- 519 Bonjean, F., and G. S. E. Lagerloef (2002), Diagnostic model and analysis of the surface currents
520 in the tropical Pacific Ocean, *J. Phys. Oceanogr.*, 32(10), 2938–2954, doi:10.1175/1520-
521 0485(2002)032<2938:DMAAOT>2.0.CO;2.
- 522 Caldeira, R. M. A., and P. Marchesiello (2002), Ocean response to wind sheltering in the
523 Southern California Bight, *Geophys. Res. Lett.*, 29(13), 1–4, doi:10.1029/2001GL014563.
- 524 Caldeira, R. M. A., S. Groom, P. Miller, D. Pilgrim, and N. P. Nezlin (2002), Sea-surface
525 signatures of the island mass effect phenomena around Madeira Island, Northeast Atlantic,
526 *Remote Sens. Environ.*, 80(2), 336–360, doi:10.1016/S0034-4257(01)00316-9.
- 527 Calil, P. H. R., K. J. Richards, Y. Jia, and R. R. Bidigare (2008), Eddy activity in the lee of the
528 Hawaiian Islands, *Deep. Res. Part II Top. Stud. Oceanogr.*, 55(10–13), 1179–1194,
529 doi:10.1016/j.dsr2.2008.01.008.
- 530 Chang, M. H., T. Y. Tang, C. R. Ho, and S. Y. Chao (2013), Kuroshio-induced wake in the lee of
531 Green Island off Taiwan, *J. Geophys. Res. Ocean.*, 118(3), 1508–1519,
532 doi:10.1002/jgrc.20151.
- 533 Chavanne, C., P. Flament, R. Lumpkin, B. Dousset, and A. Bentamy (2002), Scatterometer
534 observations of wind variations induced by oceanic islands: implications for wind-driven
535 ocean circulation., *Can. J. Remote Sens.*, 28(3), 466–474, doi:10.5589/m02-047.
- 536 Chelton, D. B., R. a. deSzoeke, M. G. Schlax, K. El Naggar, and N. Siwertz (1998),
537 Geographical Variability of the First Baroclinic Rossby Radius of Deformation, *J. Phys.*
538 *Oceanogr.*, 28(3), 433–460, doi:10.1175/1520-0485(1998)028<0433:GVOTFB>2.0.CO;2.
- 539 Chen, R., A. F. Thompson, and G. R. Flierl (2016), Time-dependent eddy-mean energy diagrams
540 and their application to the ocean, *J. Phys. Oceanogr.*, 46(9), 2827–2850, doi:10.1175/JPO-
541 D-16-0012.1.
- 542 Cummings, J. A., and O. M. Smedstad (2013), *Variational Data Assimilation for the Global*
543 *Ocean*.
- 544 Debreu, L., and E. Blayo (2008), Two-way embedding algorithms: A review, *Ocean Dyn.*, 58(5–
545 6), 415–428, doi:10.1007/s10236-008-0150-9.
- 546 Debreu, L., P. Marchesiello, P. Penven, and G. Cambon (2012), Two-way nesting in split-
547 explicit ocean models: Algorithms, implementation and validation, *Ocean Model.*, 49–50,
548 1–21, doi:10.1016/j.ocemod.2012.03.003.
- 549 Dietrich, D., M. Bowman, and C. Lin (1996), Numerical studies of small island wakes in the
550 ocean, *Geophys. Astrophys. Fluid Dyn.*, 83(3–4), 195–231,
551 doi:10.1080/03091929608208966.
- 552 DiGiacomo, P. M., and B. Holt (2001), Satellite observations of small coastal ocean eddies in the
553 Southern California Bight, *J. Geophys. Res.*, 106(10), 22,521–522,543,

- 554 doi:10.1029/2000JC000728.
- 555 Doglioli, A. M., B. Blanke, S. Speich, and G. Lapeyre (2007), Tracking coherent structures in a
556 regional ocean model with wavelet analysis: Application to Cape Basin eddies, *J. Geophys.*
557 *Res. Ocean.*, *112*(5), 1–12, doi:10.1029/2006JC003952.
- 558 Dong, C., J. C. McWilliams, and A. F. Shchepetkin (2007), Island Wakes in Deep Water, *J.*
559 *Phys. Oceanogr.*, *37*(4), 962–981, doi:10.1175/JPO3047.1.
- 560 Dong, C., X. Lin, Y. Liu, F. Nencioli, Y. Chao, Y. Guan, D. Chen, T. Dickey, and J. C.
561 McWilliams (2012), Three-dimensional oceanic eddy analysis in the Southern California
562 Bight from a numerical product, *J. Geophys. Res. Ocean.*, *117*(1), 1–17,
563 doi:10.1029/2011JC007354.
- 564 Doty, M. S., & Oguri, M. (1956), The island mass effect, *J. du Cons.*, *22*(1), 33–37,
565 doi:10.1093/icesjms/22.1.33.
- 566 Emery, W. J. (2003), Water Types and Water Masses Descriptive Tools: The TS Curve, in
567 *Ocean Circulation*, pp. 1556–1567.
- 568 Gaillard, F., T. Reynaud, V. Thierry, N. Kolodziejczyk, and K. Von Schuckmann (2016), In situ-
569 based reanalysis of the global ocean temperature and salinity with ISAS: Variability of the
570 heat content and steric height, *J. Clim.*, *29*(4), 1305–1323, doi:10.1175/JCLI-D-15-0028.1.
- 571 Halo, I., P. Penven, B. Backeberg, I. Ansorge, F. Shillington, and R. Roman (2014), Mesoscale
572 eddy variability in the southern extension of the East Madagascar Current: Seasonal cycle,
573 energy conversion terms, and eddymean properties, *J. Geophys. Res. Ocean.*, *0*, 7324–7356,
574 doi:10.1002/2014JC009820.
- 575 Harrison, D. E., and A. R. Robinson (1978), Energy analysis of open regions of turbulent flows -
576 mean eddy energetics of a numerical ocean circulation experiment, *Dyn. Atmos. Ocean.*,
577 *2*(2), 185–211, doi:10.1016/0377-0265(78)90009-X.
- 578 Hasegawa, D., H. Yamazaki, R. G. Lueck, and L. Seuront (2004), How islands stir and fertilize
579 the upper ocean, *Geophys. Res. Lett.*, *31*(16), 1–4, doi:10.1029/2004GL020143.
- 580 Hasegawa, D., H. Yamazaki, T. Ishimaru, H. Nagashima, and Y. Koike (2008), Apparent
581 phytoplankton bloom due to island mass effect, *J. Mar. Syst.*, *69*(3–4), 238–246,
582 doi:10.1016/j.jmarsys.2006.04.019.
- 583 Hasegawa, D., M. R. Lewis, and a. Gangopadhyay (2009), How islands cause phytoplankton to
584 bloom in their wakes, *Geophys. Res. Lett.*, *36*(20), 1–4, doi:10.1029/2009GL039743.
- 585 Heywood, K. J., E. D. Barton, and J. H. Simpson (1990), The effects of flow disturbance by an
586 oceanic island, *J. Mar. Res.*, *48*, 55–73, doi:10.1357/002224090784984623.
- 587 Heywood, K. J., D. P. Stevens, and G. R. Bigg (1996), Eddy formation behind the tropical island
588 of Aldabra, *Deep. Res. Part I Oceanogr. Res. Pap.*, *43*(4), 555–578, doi:10.1016/0967-
589 0637(96)00097-0.
- 590 Hristova, H. G., W. S. Kessler, J. C. McWilliams, and M. J. Molemaker (2014), Mesoscale
591 variability and its seasonality in the Solomon and Coral Seas, *J. Geophys. Res. Ocean.*,
592 *119*, 4669–4687, doi:10.1002/2013JC009741.
- 593 Jiménez, B., P. Sangrà, and E. Mason (2008), A numerical study of the relative importance of
594 wind and topographic forcing on oceanic eddy shedding by tall, deep water islands, *Ocean*
595 *Model.*, *22*(3–4), 146–157, doi:10.1016/j.ocemod.2008.02.004.
- 596 Kang, D., and E. N. Curchitser (2015), Energetics of Eddy-Mean Flow Interactions in the Gulf
597 Stream Region, *J. Phys. Oceanogr.*, *45*, 1103–1120, doi:10.1175/JPO-D-14-0200.1.
- 598 Karnauskas, K. B., G. C. Johnson, and R. Murtugudde (2017), On the climate impacts of atolls in
599 the central equatorial Pacific, *Int. J. Climatol.*, *37*(1), 197–203, doi:10.1002/joc.4697.

- 600 Kersalé, M., a. M. Doglioli, and a. a. Petrenko (2011), Sensitivity study of the generation of
601 mesoscale eddies in a numerical model of Hawaii islands, *Ocean Sci.*, 7(3), 277–291,
602 doi:10.5194/os-7-277-2011.
- 603 Large, W. G., J. C. McWilliams, and S. C. Doney (1994), Oceanic vertical mixing: A review and
604 a model with a nonlocal boundary layer parameterization, *Rev. Geophys.*, 32(4), 363,
605 doi:10.1029/94RG01872.
- 606 Legeckis, R., C. W. Brown, F. Bonjean, and E. S. Johnson (2004), The influence of tropical
607 instability waves on phytoplankton blooms in the wake of the Marquesas Islands during
608 1998 and on the currents observed during the drift of the Kon-Tiki in 1947, *Geophys. Res.
609 Lett.*, 31(23), 1–4, doi:10.1029/2004GL021637.
- 610 Liang, J. H., J. C. McWilliams, J. Kurian, F. Colas, P. Wang, and Y. Uchiyama (2012),
611 Mesoscale variability in the northeastern tropical Pacific: Forcing mechanisms and eddy
612 properties, *J. Geophys. Res. Ocean.*, 117(7), 1–13, doi:10.1029/2012JC008008.
- 613 Liu, Y., C. Dong, Y. Guan, D. Chen, J. McWilliams, and F. Nencioli (2012), Eddy analysis in
614 the subtropical zonal band of the North Pacific Ocean, *Deep. Res. Part I Oceanogr. Res.
615 Pap.*, 68, 54–67, doi:10.1016/j.dsr.2012.06.001.
- 616 Locarnini, R. A. et al. (2013), World Ocean Atlas 2013. Vol. 1: Temperature., *S. Levitus, Ed.; A.
617 Mishonov, Tech. Ed.; NOAA Atlas NESDIS*, 73(September), 40, doi:10.1182/blood-2011-
618 06-357442.
- 619 Lungu, T., and P. S. Callahan (2006), QuikSCAT science data product user’s manual: Overview
620 and geophysical data products, *D-18053-Rev A, version*, 3, 91.
- 621 Marchesiello, P., J. C. McWilliams, and A. Shchepetkin (2001), Open boundary conditions for
622 long-term integration of regional oceanic models, *Ocean Model.*, 3(1–2), 1–20,
623 doi:10.1016/S1463-5003(00)00013-5.
- 624 Marchesiello, P., J. C. McWilliams, A. Shchepetkin, P. Physics, and L. Angeles (2003),
625 Equilibrium Structure and Dynamics of the California Current System, *J. Phys. Oceanogr.*,
626 33(4), 753–783, doi:10.1175/1520-0485(2003)33<753:ESADOT>2.0.CO;2.
- 627 Martinez, E., and K. Maamaatuaiahutapu (2004), Island mass effect in the Marquesas Islands:
628 Time variation, *Geophys. Res. Lett.*, 31(18), 1–4, doi:10.1029/2004GL020682.
- 629 Martinez, E., A. Ganachaud, J. Lefevre, and K. Maamaatuaiahutapu (2009), Central south pacific
630 thermocline water circulation from a high-resolution ocean model validated against satellite
631 data: seasonal variability and el niño 1997-1998 influence, *J. Geophys. Res. Ocean.*, 114(5),
632 1–16, doi:10.1029/2008JC004824.
- 633 Martinez, E., M. Rodier, and K. Maamaatuaiahutapu (2016), Environnement océanique des
634 Marquises, , 1–15.
- 635 Mkhinini, N., A., L. S. Coimbra, A. Stegner, T. Arsouze, I. Taupier-Letage, and K. Béranger
636 (2014), Long-lived mesoscale eddies in the easternMediterranean Sea: Analysis of 20 years
637 of AVISO geostrophic velocities, *J. Geophys. Res. Ocean.*, 119, 8603–8626,
638 doi:10.1002/2014JC010176.
- 639 Neill, S. P., and A. J. Elliott (2004), Observations and simulations of an unsteady island wake in
640 the Firth of Forth, Scotland, *Ocean Dyn.*, 54(3–4), 324–332, doi:10.1007/s10236-003-0084-
641 1.
- 642 Nencioli, F., C. Dong, T. Dickey, L. Washburn, and J. C. Mcwilliams (2010), A vector
643 geometry-based eddy detection algorithm and its application to a high-resolution numerical
644 model product and high-frequency radar surface velocities in the Southern California Bight,
645 *J. Atmos. Ocean. Technol.*, 27(3), 564–579, doi:10.1175/2009JTECHO725.1.

- 646 O'Connor, B. M., R. A. Fine, and D. B. Olson (2005), A global comparison of subtropical
647 underwater formation rates, *Deep. Res. Part I Oceanogr. Res. Pap.*, 52(9), 1569–1590,
648 doi:10.1016/j.dsr.2005.01.011.
- 649 Palacios, D. M. (2002), Factors influencing the island-mass effect of the Galápagos Archipelago,
650 *Geophys. Res. Lett.*, 29(23), 1–4, doi:10.1029/2002GL016232.
- 651 Penven, P., P. Marchesiello, L. Debreu, and J. Lefèvre (2007), Software tools for pre- and post-
652 processing of oceanic regional simulations, *Environ. Model. Softw.*, 23(5), 660–662,
653 doi:10.1016/j.envsoft.2007.07.004.
- 654 Qiu, B., and S. Chen (2004), Seasonal Modulations in the Eddy Field of the South Pacific Ocean,
655 *J. Phys. Oceanogr.*, 34(7), 1515–1527, doi:10.1175/1520-
656 0485(2004)034<1515:SMITEF>2.0.CO;2.
- 657 Qu, T., S. Gao, and R. A. Fine (2013), Subduction of South Pacific Tropical Water and Its
658 Equatorward Pathways as Shown by a Simulated Passive Tracer*, *J. Phys. Oceanogr.*,
659 43(8), 1551–1565, doi:10.1175/JPO-D-12-0180.1.
- 660 Sangrà, P., G. Basterretxea, J. L. Pelegrí, and J. Aristegui (2001), Chlorophyll increase due to
661 internal waves on the shelf break of Gran Canaria (Canary Islands)*, *Signals*, 65, 89–97,
662 doi:http://digital.csic.es/bitstream/10261/5326/1/wave.pdf.
- 663 Sangrà, P., J. L. Pelegrí, A. Hernández-Guerra, I. Arregui, J. M. Martín, A. Marrero-Díaz, A.
664 Martínez, A. W. Ratsimandresy, and A. Rodríguez-Santana (2005), Life history of an
665 anticyclonic eddy, *J. Geophys. Res. C Ocean.*, 110(3), 1–19, doi:10.1029/2004JC002526.
- 666 Shchepetkin, A. F., and J. C. McWilliams (2003), A method for computing horizontal pressure-
667 gradient force in an oceanic model with nonaligned vertical coordinate, *J. Geophys. Res.*,
668 108(C3), 1–34, doi:10.1029/2001JC001047.
- 669 Shchepetkin, A. F., and J. C. McWilliams (2005), The regional oceanic modeling system
670 ({ROMS}): a split-explicit, free-surface, topography-following-coordinate oceanic model,
671 *Ocean Model.*, 9, 347–404, doi:10.1016/j.ocemod.2004.08.002.
- 672 Signorini, S. R., C. R. McClain, and Y. Dandonneau (1999), Mixing and phytoplankton bloom in
673 the wake of the Marquesas Islands, *Geophys. Res. Lett.*, 26(20), 3121–3124,
674 doi:10.1029/1999GL010470.
- 675 da Silva, A. M., C. C. Young, and S. Levitus (1994), Atlas of surface marine data 1994, vol. 4,
676 Anomalies of fresh water fluxes, *Noaa atlas nesdis*, 9, 308.
- 677 Smith, W. H., and D. Sandwell (1997), Global Sea Floor Topography from Satellite Altimetry
678 and Ship Depth Soundings, *Science (80-.)*, 277(5334), 1956–1962,
679 doi:10.1126/science.277.5334.1956.
- 680 Stegner, A. (2014), Oceanic Island Wake Flows in the Laboratory, in *Modeling Atmospheric and*
681 *Oceanic Flows: Insights from Laboratory Experiments and Numerical Simulations*, pp.
682 265–276.
- 683 Sun, Z., Z. Zhang, W. Zhao, and J. Tian (2016), Interannual modulation of eddy kinetic energy
684 in the northeastern South China Sea as revealed by an eddy-resolving OGCM, *J. Geophys.*
685 *Res. Ocean.*, 121, 1–12, doi:10.1002/2015JC011497.
- 686 Tomczak, M. (1988), Island wakes in deep and shallow water, *J. Geophys. Res.*, 93(8), 5153–
687 5154.
- 688 Wolanksi, E., T. Asaeda, A. Tanaka, and E. Deleersnijder (1996), Three-dimensional island
689 wakes in the field, laboratory experiments and numerical models, *Cont. Shelf Res.*, 16(11),
690 1437–1452, doi:10.1016/0278-4343(95)00087-9.
- 691 Wyrski, K. (1981), An Estimate of Equatorial Upwelling in the Pacific, *J. Phys. Oceanogr.*, 11,

692 1205–1214.

693 Xu, Y., and R. B. Scott (2008), Subtleties in forcing eddy resolving ocean models with satellite
694 wind data, *Ocean Model.*, 20(3), 240–251, doi:10.1016/j.ocemod.2007.09.003.

695 Yan, X., C. Ho, Q. Zheng, and V. Klemas (1992), Temperature and Size Variabilities of the
696 Western Pacific Warm Pool, *Science (80-.)*, 258(5088), 1643–1645.

697 Zweng, M. M., J. R. Reagan, J. I. Antonov, A. V. Mishonov, T. P. Boyer, H. E. Garcia, O. K.
698 Baranova, D. R. Johnson, D. Seidov, and M. M. Bidlle (2013), World Ocean Atlas 2013,
699 Volume 2: Salinity, *NOAA Atlas NESDIS, 119(1)*, 227–237, doi:10.1182/blood-2011-06-
700 357442.

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702

703 Figure 1: (a) Mean surface current from the satellite derived Ocean Surface Current Analysis –
704 Real time (OSCAR) product (time averaged over Oct-1992 to Jun-2015, in m s^{-1}) over French
705 Polynesia. The islands are represented in white, as well as the name of the five archipelagos. The
706 black boxes show the parent and child grids implemented in the ROMS configuration. (b)
707 Bottom topography from the 2-arc minute topography/bathymetry dataset ETOPO2 (in m)
708 around the archipelago used in the model configuration as well as the names of the main islands.

709

710 Figure 2: SST ($^{\circ}\text{C}$) from Geostationary Operational Environmental Satellite (GOES)/ Polar
711 Operational Environmental Satellite (POES) in the Marquesas archipelago for (a) 2015-NOV-17
712 and (b) 2012-AUG-21.

713

714 Figure 3: Mean SST over the child grid area from (a) MODIS (time averaged from Jul-2002 to
715 Jun-2015) and (b) W13Q at 10m (time averaged over year 4 to 10). (c) Monthly climatology
716 averaged over the child grid area from MODIS (black), ISAS13 (red) and W13Q (green). The
717 grey shaded area represents the spatial standard deviation from MODIS monthly climatology.
718 Units are in $^{\circ}\text{C}$.

719

720 Figure 4: (a) Time series of the monthly averaged salinity (in psu) at 10 m issued from ISAS13
721 (red) and W13Q (green). Red shaded area corresponds to the spatial standard deviation from
722 ISAS13. (b) TS diagram from ISAS13 (red) and W13Q (green) monthly climatology and their
723 standard deviations (red and green shades respectively).

724

725 Figure 5: (a) W13Q current at 10 m time averaged over year 4 to 10 and over the child domain.
726 For clarity, only 1 vector each 17 is represented. (b) Zoom of (a) over the Marquesas northern
727 islands, as defined by the black box on a). For clarity, only 1 vector each 3 is represented. Units
728 are in m s^{-1} .

729

730 Figure 6: Annual mean of EKE_m (in $\text{cm}^2 \text{s}^{-2}$) obtained from monthly averaged of (a) OSCAR
731 surface currents and (b) W13Q currents at 10m. (c) is the annual mean of EKE derived from the
732 W13Q 2-day outputs. Their standard deviations are represented in (d), (e) and (f), respectively.

733

734 Figure 7: Spatial distribution of the annual mean of the (a) barotropic ($KmKe$), (b) baroclinic
735 ($PeKe$) energy conversion both integrated over the surface to 100 m and (c) the EKE generation
736 due to transient wind, from year 4 to 10 of W13Q (units are in $\text{cm}^3 \text{s}^{-3}$). (d) Time series of wind
737 work, baroclinic and barotropic energy (left axis), in yellow, blue and red, respectively. The sum

738 of these three terms is indicated in black. The *EKE* (in $\text{cm}^2 \text{s}^{-2}$, right axis) spatially averaged over
739 the child grid and at 10 m is represented in green.

740

741 Figure 8: Snapshots of relative vorticity at 10 m (units are in s^{-1}) from W13Q during Year 9
742 every 6 days (a-f). The detected eddies are represented by *green* contours (Nencioli et al. [2010]
743 based algorithm). The animation of the vorticity from June to July is available in the auxiliary
744 material.

745

746 Figure 9: (a) Location of eddy generation identified from W13Q over year 4 to 10. Anticyclonic
747 and cyclonic eddies are represented by red and blue circles, respectively. (b) Monthly mean
748 number of generated eddies from W13Q from year 4 to 10 (green line) and wind stress issued
749 from QuickSCAT in N.m^{-2} (orange line).

750

751 Figure 10: (a) Distribution of the detected eddy radius from year 4 to 10 in W13Q (green) and
752 from 2006 to 2012 in HYCOM (black). (b) Distribution of W13Q cyclonic (blue) vs.
753 anticyclonic (red) eddies.

754

755 Figure 11: Eddy generation sites (black circles) and their trajectories for (a) anticyclonic and (b)
756 cyclonic eddies detected in W13Q and for a maximum radius over 20-30 km encountered over
757 their lifetime. The size of the radius along the trajectory is given by the color bar (in km).

758

759 Figure 12: Snapshots of vertical velocities (in m day^{-1}) as in Figure 8, but at the mixed layer
760 depth.

761